

# EMaII PDEs

## Example Sheet 5: Inhomogeneous equations and d'Alembert

2007

**1(a)** Let  $u = v(x, y) + w(x, y)$ , where  $w$  solves the homogeneous problem. In this case we shall

$$\text{try } v(x, y) = e^x + e^y.$$

Note that  $v_{xx} = e^x$  and  $v_{yy} = e^y$ . Hence

$$v_{xx} + v_{yy} = e^x + e^y,$$

and so  $v$  is a particular integral of the inhomogeneous problem.

Now, consider the boundary conditions. using  $u(x, y) = v(x, y) + w(x, y)$

$$\begin{aligned} 0 = u(0, y) &= 1 + e^y + w(0, y) \Rightarrow w(0, y) = -1 - e^y \\ 0 = u(a, y) &= e^a + e^y + w(a, y) \Rightarrow w(a, y) = -e^a - e^y \\ 0 = u(x, 0) &= e^x + 1 + w(x, 0) \Rightarrow w(x, 0) = -1 - e^x \\ 0 = u(x, b) &= e^x + e^b + w(x, b) \Rightarrow w(x, b) = -e^b - e^x \end{aligned}$$

**1(b)** We set  $u(x, t) = v(x) + w(x, t)$ , where we

$$\text{try } v(x) = A \sin(x).$$

Substituting this form into the PDE  $v_{tt} - c^2 v_{xx} = \sin(x)$  gives

$$-c^2 A \sin(x) = \sin(x),$$

hence  $A = -1/c^2$ . So we have

$$v(x) = -(1/c^2) \sin(x), \quad w_{tt} - c^2 w_{xx} = 0.$$

Now, let us consider the boundary conditions. We get

$$\begin{aligned} 0 = u(0, t) &= w(0, t) \Rightarrow w(0, t) = 0 \\ 0 = u(\pi, t) &= 0 + w(\pi, t) \Rightarrow w(\pi, t) = 0 \\ 0 = u(x, 0) &= -(1/c^2) \sin(x) + w(x, 0) \Rightarrow w(x, 0) = \sin(x)/c^2 \\ 0 = u_t(x, b) &= -(1/c^2) \cos(x) + w_t(x, b) \Rightarrow w_t(x, b) = \cos(x)/c^2 \end{aligned}$$

**1(c)** Let  $u(x, t) = v(x, t) + w(x, t)$  and

$$\text{try } v(x, t) = Ax^2t + Ct^2.$$

Substituting this form into the PDE we get

$$Ax^2 + 2Ct - 2A\alpha^2t = x^2.$$

The coefficient of  $x^2$  gives  $A = 1$ , and the coefficient of  $t$  gives  $2C = 2A\alpha^2$ , so  $C = \alpha^2$ . Hence, we get

$$v(x, t) = x^2t + \alpha^2t^2, \quad \text{and} \quad w_t - \alpha^2w_{xx} = 0.$$

Now, from the boundary conditions we obtain

$$\begin{aligned} 0 = u(x, 0) = 0 + w(x, 0) &\Rightarrow w(x, 0) = 0 \\ 0 = u(0, t) = \alpha^2t^2 &\Rightarrow w(0, t) = -\alpha^2t^2 \\ 0 = u(L, t) = L^2t + \alpha^2t^2 + w(L, t) &\Rightarrow w(L, t) = -L^2t - \alpha^2t^2 \end{aligned}$$

**2(a)** Let  $u(x, t)$  represent the temperature field at position  $x$  and time  $t$  within the bar. The fact that ends  $x = 0$  and  $x = L$  are held at constant temperatures  $T_1$  and  $T_2$  respectively implies

$$u(0, t) = T_1, \quad u(L, t) = T_2.$$

The initial temperature is held constant at  $T_0$ , so

$$u(x, 0) = T_0.$$

Finally, heat conduction within the bar, given a thermal condition constant  $\alpha^2$ , satisfies the heat equation

$$u_t = \alpha^2u_{xx} \quad 0 \leq x < L, \quad 0 \leq t < \infty.$$

**2(b)** To solve this equation we start by letting  $u = v(x) + w(x, t)$  where  $v$  is chosen to solve the spatial boundary values

$$v(0) = T_1, \quad v(L) = T_2.$$

Hence, we choose

$$v = T_1(1 - x/L) + T_2x/L.$$

Note that such a function satisfies the PDE, because  $v_t = 0$  by definition and  $v_{xx} = 0$  since  $v$  is only linear in  $x$ .

**2(c)** Using the principle of linear superposition  $w(x, t)$  satisfies the same PDE, but with different boundary conditions. To work out the new boundary conditions we set  $u(x, t) = v(x) + w(x, t)$  and apply to the original boundary conditions. Thus:

$$\begin{aligned} T_1 = u(0, t) = T_1 + w(0, t) &\Rightarrow w(0, t) = 0 \\ T_2 = u(L, t) = T_2 + w(L, t) &\Rightarrow w(L, t) = 0 \\ T_0 = u(x, 0) = T_1(1 - x/L) + T_2x/L &\Rightarrow w(x, 0) = (T_0 - T_1) + (x/L)(T_1 - T_2). \end{aligned}$$

**2(d)** Now, if we let  $T_0 - T_1 := a$  and  $\frac{T_1 - T_2}{L} := b$  then, by definition we have

$$\begin{aligned} b_n &= \frac{2}{L} \int_0^L (a + bx) \sin(n\pi x/L) dx \\ &= \frac{2}{L} \left( \left[ -\frac{L(ax + b)}{n\pi} \cos(n\pi x/L) \right]_0^L + \int_0^L \frac{Lb}{n\pi} \cos(n\pi x/L) dx \right) \\ &= \left[ -\frac{2(ax + b)}{n\pi} \cos(n\pi x/L) + \frac{2Lb}{n^2\pi^2} \sin(n\pi x/L) \right]_0^L \\ &= \left[ \frac{(a + bL)}{2n\pi} \cos(n\pi) \right] - \left[ \frac{2b}{n\pi} \cos(n\pi) \right] \\ &= \left[ \frac{2(a + bL)}{n\pi} (-1)^n \right] - \left[ \frac{2b}{n\pi} \right] \end{aligned}$$

Hence we have

$$u(x, t) = a + bx + \sum_{n=1}^{\infty} 2 \frac{(a + bL)\pi^2(-1)^n - b}{n\pi} \sin(n\pi xL),$$

where  $a = T_0 - T_1$  and  $b = \frac{T_1 - T_2}{L}$ .

**3(a)** We have

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0, \quad 0 \leq x < 1, \quad 0 \leq y < 1$$

$u(0, y) = 0$ ,  $u(1, y) = Cy$ ,  $u(x, 0) = 0$ ,  $u(x, 1) = Cx$ . Now, if we let  $u(x, y) = v(x, y) + w(x, y)$ , we can stipulate that both  $v$  and  $w$  should satisfy Laplace's equation, but that each one takes a different inhomogeneous boundary condition. So we can choose

$$v(0, y) = 0, \quad v(1, y) = Cy, \quad v(x, 0) = 0, \quad v(x, 1) = 0,$$

and

$$w(0, y) = 0, \quad w(1, y) = 0, \quad w(x, 0) = 0, \quad w(x, 1) = Cx.$$

**3(b)** Now we can use the separation of variables method for  $v$  and  $w$  separately. Starting with  $v(x, y)$ , we set

$$v(x, y) = X_1(x)Y_1(y).$$

Substituting this function into Laplace's equation and dividing by  $X_1Y_1$  we obtain

$$\frac{X_1''}{X_1}(x) = -\frac{Y_1''}{Y_1}(y) = k_1^2.$$

Notice that we choose the constant  $k_1^2$  to be positive because we want sinusoidal solutions in  $y$  in order to satisfy the inhomogeneous boundary condition  $= Cy$ . Thus we get the two separate ODEs

$$X_1'' = k_1 X_1, \quad Y_1'' = -k_1 Y_1,$$

which have solution

$$X_1(x) = A \cosh(k_1 x) + B \sinh(k_1 x), \quad Y(y) = C \cos(k_1 y) + D \sin(k_1 y)$$

Applying the homogeneous boundary conditions for  $v$  we have

$$\begin{aligned} 0 = v(0, y) &\Rightarrow X_1(0)Y_1(y) = A \cosh(0)Y_1(y) \Rightarrow A = 0, \\ 0 = v(x, 0) &\Rightarrow X_1(x)Y_1(0) = X_1(x)C \cos(0) \Rightarrow C = 0 \\ 0 = v(x, 1) &\Rightarrow X_1(x)Y_1(1) = X_1(x)D \sin(k_1) \Rightarrow k_1 = n\pi, \end{aligned}$$

for some integer  $n$ . Hence we have

$$v(x, y) = \sum_{n=1}^{\infty} b_n \sin(n\pi y) \sinh(n\pi x),$$

where  $CD = b_n$ . Now, solving the inhomogeneous boundary condition, we have

$$Cy = v(1, y) = \sum_{n=1}^{\infty} b_n \sin(n\pi y) \sinh(n\pi),$$

Hence, from the Fourier half-range sine series we have

$$b_n \sinh(n\pi) = 2 \int_0^1 Cy \sin(n\pi y) dy$$

Solving this integral, we get

$$\begin{aligned}
 b_n \sinh(n\pi) &= 2 \int_0^1 C y \sin(n\pi y) dy \\
 &= 2 \left( \left[ -\frac{C y}{n\pi} \cos(n\pi y) \right]_0^1 + \int_0^1 \frac{C}{n^2 \pi^2} \cos(n\pi y) dy \right) \\
 &= \left[ -\frac{2C y}{n\pi} \cos(n\pi y) + \frac{2C}{n^2 \pi^2} \sin(n\pi y) \right]_0^1 \\
 &= \left[ -\frac{2C}{n\pi} \cos(n\pi) \right] - [0] \\
 &= -\frac{2C(-1)^{n+1}}{n\pi}
 \end{aligned}$$

Similarly, we can solve for  $w(x, y)$ . Note that  $w$  satisfies exactly the same PDE as  $v$  with the same boundary conditions BUT with  $x$  replaced by  $y$  and vice versa. Hence we have

$$w(x, y) = \sum_{n=1}^{\infty} b_n \sin(n\pi x) \sinh(n\pi y),$$

where the  $b_n$  are exactly as above. Hence we get the general solution  $u = v + w$ :

$$u(x, y) = \sum_{n=1}^{\infty} \frac{2C(-1)^{n+1}}{n\pi \sinh(n\pi)} (\sin(n\pi x) \sinh(n\pi y) + \sinh(n\pi x) \sin(n\pi y))$$

**4(a)** We start by stating the d'Alembert solution

$$u(x, t) = f(x - ct) + g(x + ct).$$

Posing the initial condition  $u_t(x, 0) = 0$ , we find

$$0 = -cf'(x) + cg'(x), \quad \Rightarrow g(x) - f(x) = K.$$

for some constant  $K$ . The other initial condition,  $u(x, 0) = e^{-x^2} \sin(x)$ , gives

$$f(x) + g(x) = e^{-x^2} \sin(x)$$

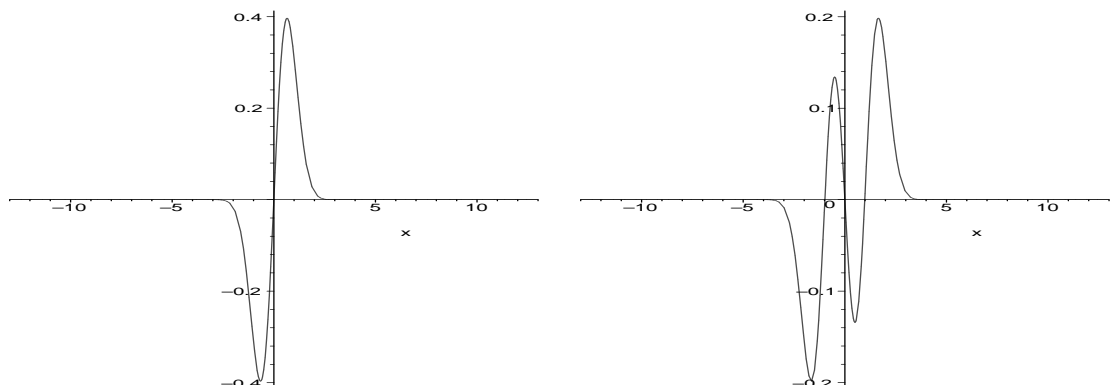
Hence we have

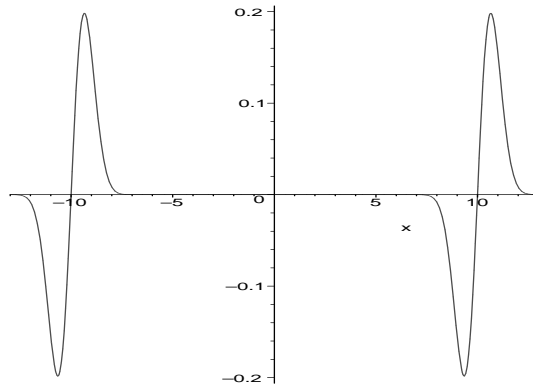
$$\begin{aligned}
 g(x) &= \frac{K}{2} + \frac{1}{2} e^{-x^2} \sin(x) \\
 f(x) &= -\frac{K}{2} + \frac{1}{2} e^{-x^2} \sin(x).
 \end{aligned}$$

Which gives

$$u(x, t) = f(x - ct) + g(x + ct) = \frac{1}{2} \left( e^{-(x+ct)^2} \sin(x + ct) - e^{-(x-ct)^2} \sin(x - ct) \right).$$

The solutions for  $t = 0$ ,  $t = 1$  and  $t = 10$  with  $c = 1$  are plotted below





4(b) Again we write down d'Alembert's solution

$$u(x, t) = f(x - ct) + g(x + ct)$$

and pose the boundary conditions. The initial velocity condition gives

$$e^{-|x|} \cos(x) = -cf'(x) + cg'(x).$$

Which implies that

$$g(x) - f(x) = \frac{1}{c} \int e^{-|x|} \cos(x) dx$$

Let  $I = \int e^{-|x|} \cos(x) dx$ . Then, integration by parts twice yields

$$\begin{aligned} I &= [-e^{-|x|} \sin(x)] + \int \text{sign}(x) e^{-|x|} \sin(x) dx \\ &= [-e^{-|x|} \sin(x)] + [\text{sign}(x) e^{-|x|} \cos(x)] - \int (\text{sign}(x))^2 e^{-|x|} \cos(x) dx \\ &= [-e^{-|x|} \sin(x) + \text{sign}(x) e^{-|x|} \cos(x)] - I. \end{aligned}$$

Hence,

$$I = \frac{1}{2}(\text{sign}(x) e^{-|x|} \cos(x) - e^{-|x|} \sin(x)) + K$$

and so

$$g(x) - f(x) = \frac{1}{2c}(\text{sign}(x) e^{-|x|} \cos(x) - e^{-|x|} \sin(x)) + K \quad (1)$$

The initial condition on position gives

$$u(x, 0) = f(x) + g(x) = \frac{1}{1 + x^2} \quad (2)$$

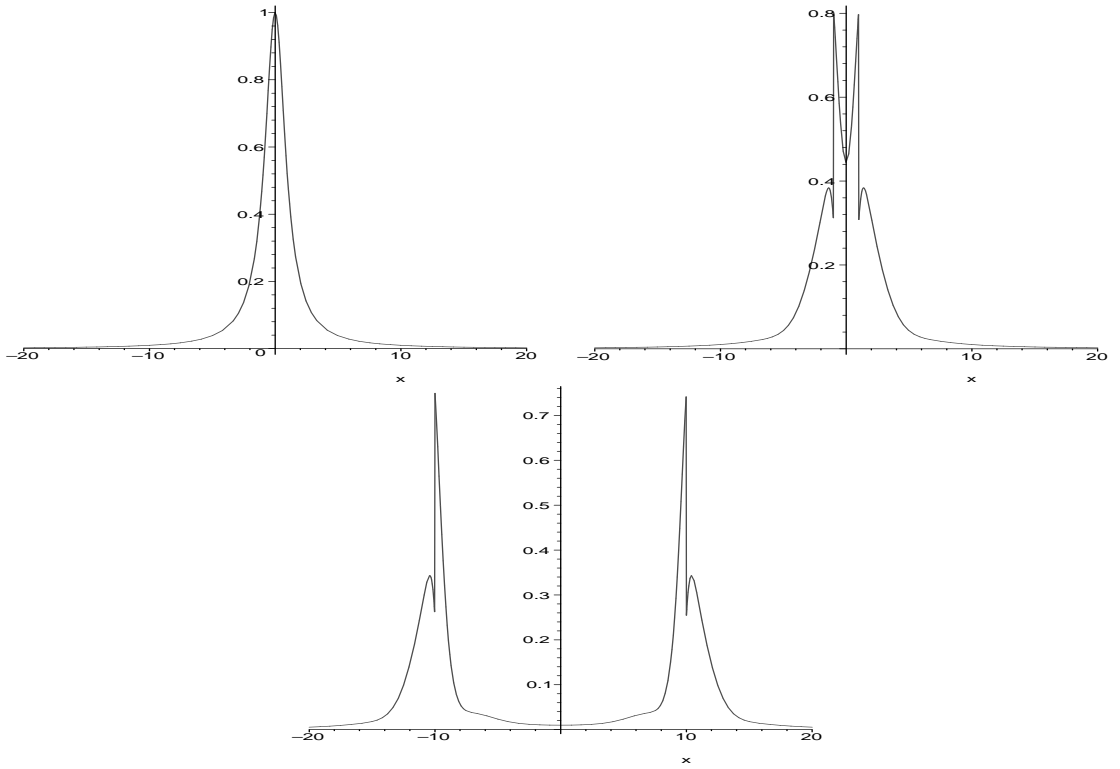
Solving the simultaneous equations (1) and (2) gives

$$\begin{aligned} g(x) &= \frac{1}{2(1 + x^2)} + \frac{1}{4c}(\text{sign}(x) e^{-|x|} \cos(x) - e^{-|x|} \sin(x)) + K/2 \\ f(x) &= \frac{1}{2(1 + x^2)} - \frac{1}{4c}(\text{sign}(x) e^{-|x|} \cos(x) - e^{-|x|} \sin(x)) - K/2. \end{aligned}$$

Hence  $u(x, t) = f(x - ct) + g(x + ct)$  gives

$$\begin{aligned} u(x, t) &= \frac{1}{2[1 + (x - ct)^2]} + \frac{1}{2(x + ct)^2} + \frac{1}{4c}(\text{sign}(x + ct) e^{-|x+ct|} \cos(x + ct) - e^{-|x+ct|} \sin(x + ct)) \\ &\quad - \frac{1}{4c}(\text{sign}(x - ct) e^{-|x-ct|} \cos(x - ct) - e^{-|x-ct|} \sin(x - ct)) \end{aligned}$$

The solutions for  $t = 0$ ,  $t = 1$  and  $t = 10$  with  $c = 1$  are plotted below



5(a) We start by writing down the d'Alembert solution

$$u(x, t) = f(x - ct) + g(x + ct).$$

Posing the initial velocity condition we get

$$e^{-x} = u_t(x, 0) = -cf'(x) + cg'(x).$$

Hence,

$$g(x) - f(x) = -\frac{1}{c}e^{-x} + K, \tag{3}$$

for some arbitrary constant  $K$ . Posing the initial condition on position we get

$$\frac{1}{x^2} = u(x, 0) = f(x) + g(x).$$

Solving the simultaneous equations (3) and () we find

$$\begin{aligned} f(x) &= \frac{1}{2(1+x)^2} + \frac{1}{2c}e^{-x} - \frac{K}{2} \\ g(x) &= \frac{1}{2(1+x)^2} - \frac{1}{2c}e^{-x} + \frac{K}{2} \end{aligned}$$

both of which are only valid for  $x > 0$ . Hence we find

$$g(x + ct) = \frac{1}{2(1+x+ct)^2} - \frac{1}{2c}e^{-x-ct} + \frac{K}{2}$$

for all  $x + ct > 0$  (which is true for all  $t > 0$ ), and

$$f(x - ct) = \frac{1}{2(1+x-ct)^2} + \frac{1}{2c}e^{-x+ct} - \frac{K}{2}$$

for  $x - ct > 0$ .

Hence, it remains to solve for  $f(x - ct)$  for  $x < ct$ . Do do this we use the boundary condition at  $t = 0$ . Hence we obtain

$$\begin{aligned}\sin(\pi t) &= u(0, t) = f(-ct) + g(ct) \\ &= f(-ct) + \frac{1}{2(1 + ct)^2} - \frac{1}{2c}e^{-ct} + \frac{K}{2},\end{aligned}$$

from which we note

$$f(-ct) = \cos(\pi t) - \frac{1}{2(1 + ct)^2} + \frac{1}{2c}e^{-ct} - \frac{K}{2},$$

for  $t > 0$ . Letting  $z = -ct$ , we find

$$f(z) = \cos\left(\frac{\pi z}{c}\right) - \frac{1}{2(1 - z)^2} + \frac{1}{2c}e^z - \frac{K}{2},$$

for an arbitrary argument  $z < 0$ . Hence we obtain

$$f(x - ct) = \cos\left(\frac{\pi x - ct}{c}\right) - \frac{1}{2(1 - x + ct)^2} + \frac{1}{2c}e^{x-ct} - \frac{K}{2},$$

for  $x - ct < 0$ .

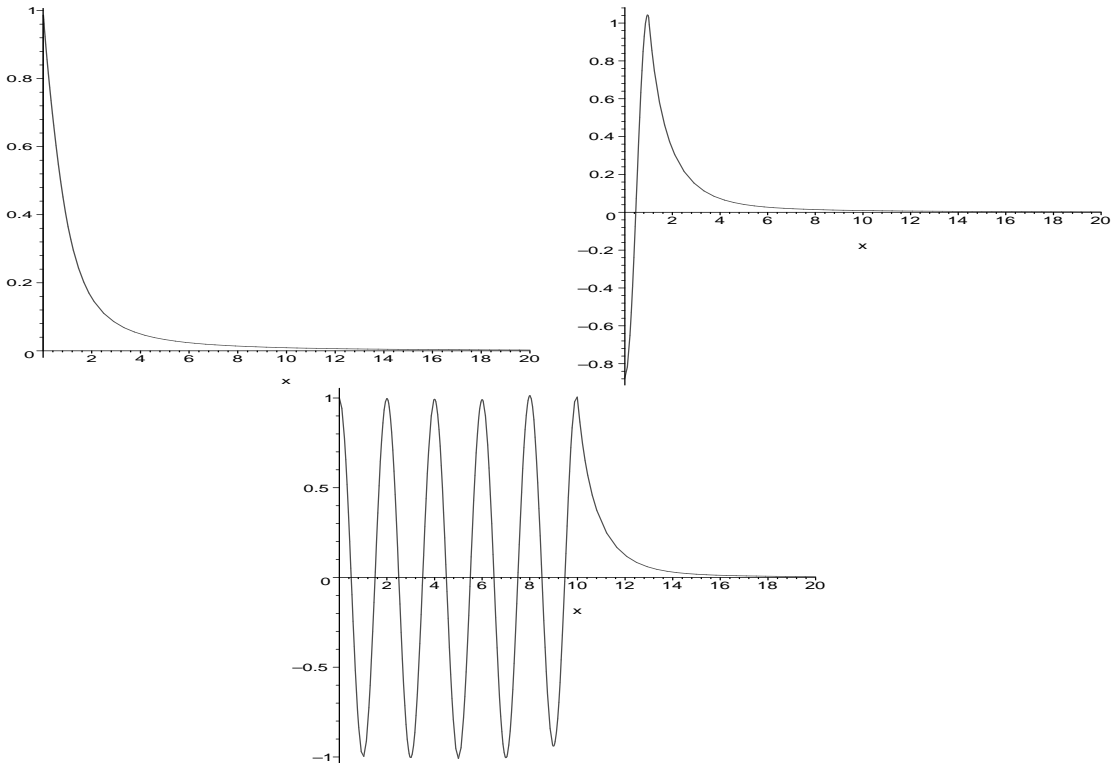
Adding  $f$  and  $g$ , we obtain

$$u(x, t) = \frac{1}{2(1 + x + ct)^2} - \frac{1}{2c}e^{-x-ct} + \frac{1}{2(1 + x - ct)^2} + \frac{1}{2c}e^{-x+ct}$$

for  $x > ct$  and

$$u(x, t) = \frac{1}{2(1 + x + ct)^2} - \frac{1}{2c}e^{-x-ct} \cos\left(\frac{\pi x - ct}{c}\right) - \frac{1}{2(1 - x + ct)^2} + \frac{1}{2c}e^{x-ct}$$

for  $x < ct$ . The solution for  $ct = 0$ ,  $ct = 1$  and  $ct = 10$  is plotted below



5(b) We start by writing down the d'Alembert solution

$$u(x, t) = f(x - ct) + g(x + ct).$$

Posing the initial velocity condition we get

$$e^{-x} = u_t(x, 0) = -cf'(x) + cg'(x).$$

Hence,

$$g(x) - f(x) = -\frac{1}{c}e^{-x} + K, \quad (4)$$

for some arbitrary constant  $K$ . Posing the initial condition on position we get

$$0 = u(x, 0) = f(x) + g(x).$$

Solving the simultaneous equations (4) and () we find

$$\begin{aligned} f(x) &= \frac{1}{2c}e^{-x} - \frac{K}{2}, \\ g(x) &= -\frac{1}{2c}e^{-x} + \frac{K}{2} \end{aligned}$$

both of which are only valid for  $x > 0$ . Hence we find

$$g(x + ct) = -\frac{1}{2c}e^{-x-ct} + \frac{K}{2},$$

for all  $x + ct > 0$  (which is true for all  $t > 0$ ), and

$$f(x - ct) = \frac{1}{2c}e^{-x+ct} - \frac{K}{2},$$

for  $x - ct > 0$ .

It remains to solve for  $f(x - ct)$  for  $x < ct$ . Do do this we use the boundary condition at  $t = 0$ . Thus, we obtain

$$\begin{aligned} \frac{1}{2} - \cos^2(2\pi t) &= u_x(0, t) = f'(-ct) + g'(ct) \\ &= f'(-ct) - \frac{1}{2c}e^{-ct}, \end{aligned}$$

from which we get

$$f'(-ct) = \frac{1}{2} - \cos^2(2\pi t) + \frac{1}{2c}e^{-ct}$$

Letting  $z = -ct$  we find

$$f'(z) = \frac{1}{2} - \cos^2\left(\frac{2\pi z}{c}\right) + \frac{1}{2c}e^z,$$

for any argument  $z < 0$ . Upon integration we find

$$f(z) = -\frac{c}{8\pi} \sin\left(\frac{4\pi z}{c}\right) + \frac{1}{2c}e^z + \text{const.}$$

Hence

$$f(x - ct) = -\frac{c}{8\pi} \sin\left(\frac{4\pi(x - ct)}{c}\right) + \frac{1}{2c}e^{x-ct} + \text{const..}$$

We choose the constant equal to  $-K/2$  in order to cancel the  $+K/2$  in  $g(x + ct)$ .

We get the general solution upon adding  $f(x - ct)$  and  $g(x + ct)$  for the two different cases.

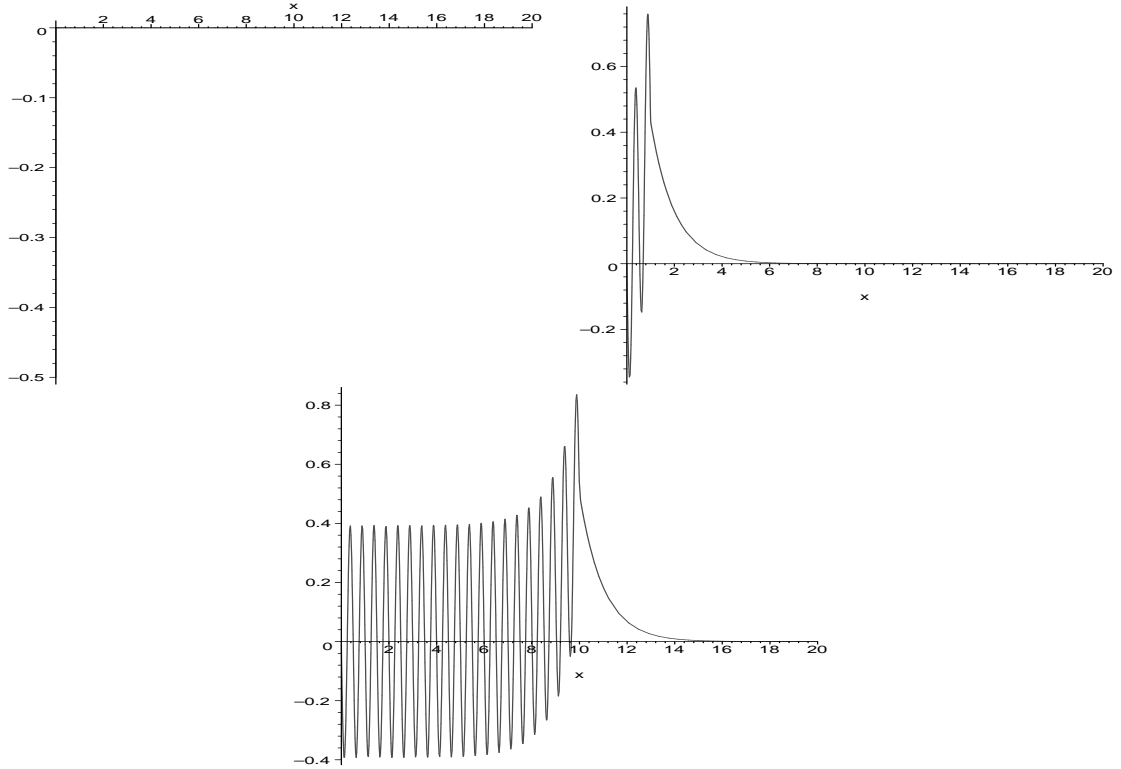
For  $x < ct$  we have:

$$u(x, t) = -\frac{c}{8\pi} \sin\left(\frac{4\pi(x - ct)}{c}\right) + \frac{1}{2c}e^{x-ct} - \frac{1}{2c}e^{-x-ct}$$

and for  $x > ct$ :

$$u(x, t) = +\frac{1}{2c}e^{-x+ct} - \frac{1}{2c}e^{-x-ct}$$

The solution for  $ct = 0$ ,  $ct = 1$  and  $ct = 10$  is plotted below



Note that in the first graph the solution is zero everywhere.

**6** The question leads you through the solution to this. If not, look at the relevant page of Kreyszig.